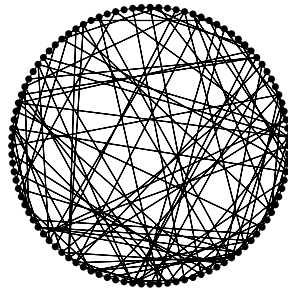


Disordered (Soft) Spin Systems on Random Graphs with Finite Connectivity

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Disordered Many-Particle Systems

Finite Connectivity – Statics

- Attractor Neural Networks
- ‘Small World’ Models
- Soft Spin RS: Order Parameter Functionals

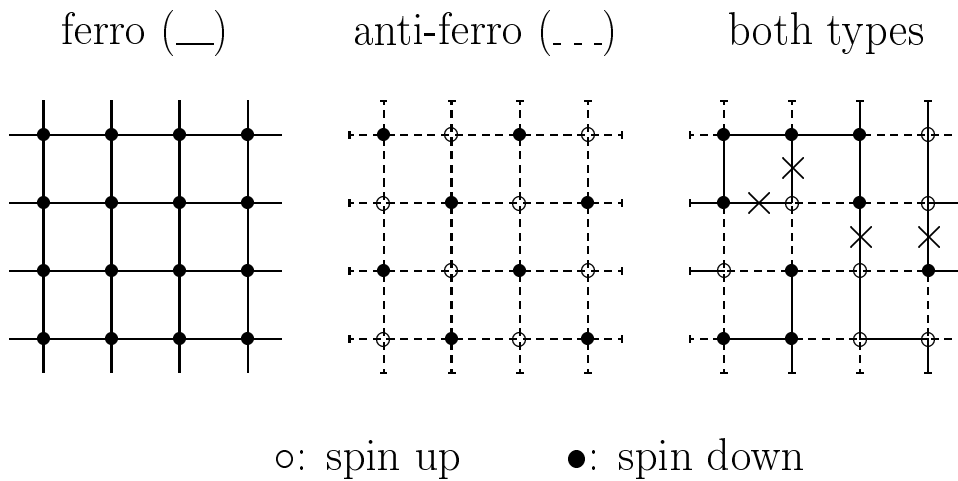
Finite Connectivity – Dynamics

- Generating Functional Analysis

Outlook

DISORDERED MANY-PARTICLE SYSTEMS

- (pseudo-) randomness in microscopic parameters
- no microscopic periodicity
- high degree of frustration, incompatible forces



- many relevant time-scales
- many techniques no longer applicable
- even mean-field models are non-trivial

physics: spin-glasses, glasses
 biology: neural networks, proteins, immune networks
 IT: machine learning, error-correcting codes
 economics: agent-based market models

The Replica Method (Marc Kac, 1968)



- $x^n = 1 + n \log x + \mathcal{O}(n^2)$

$$\overline{\log Z[\mathbf{J}]} = \lim_{n \rightarrow 0} \frac{1}{n} \log \overline{Z^n[\mathbf{J}]}$$

- n integer:

$$Z^n[\mathbf{J}] = \left[\sum_{\boldsymbol{\sigma}} e^{-\beta H(\boldsymbol{\sigma})} \right]^n = \sum_{\boldsymbol{\sigma}^1} \dots \sum_{\boldsymbol{\sigma}^n} e^{-\beta \sum_{\alpha=1}^n H(\boldsymbol{\sigma}^\alpha)}$$

partition function of n *independent* replicas of system

- disorder average:

$$\overline{f} = - \lim_{n \rightarrow 0} \frac{1}{\beta N n} \log \left[\sum_{\boldsymbol{\sigma}^1} \dots \sum_{\boldsymbol{\sigma}^n} e^{-\beta H_{\text{eff}}(\boldsymbol{\sigma}^1, \dots, \boldsymbol{\sigma}^n)} \right]$$

$$H_{\text{eff}}(\boldsymbol{\sigma}^1, \dots, \boldsymbol{\sigma}^n) = -\frac{1}{\beta} \log \overline{e^{-\beta \sum_{\alpha=1}^n H(\boldsymbol{\sigma}^\alpha)}}$$

n *interacting* replicas, *no* disorder

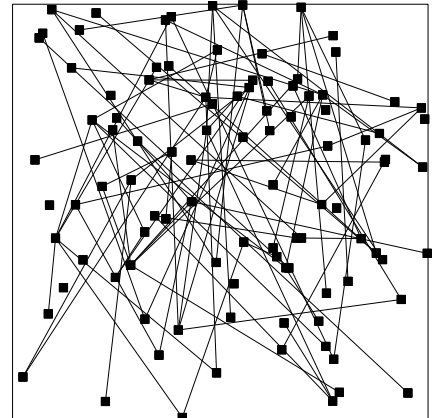
- exchange limits

$N \rightarrow \infty$ versus $n \rightarrow 0$

FINITE CONNECTIVITY STATICS

N spins on random graph, $c_{ij} \in \{0, 1\}$

$$H = - \sum_{i < j} c_{ij} J_{ij} \sigma_i \sigma_j + \sum_i V(\sigma_i)$$



$N = 100, c = 2$

- $P(c_{ij}) = \frac{c}{N} \delta_{c_{ij},1} + (1 - \frac{c}{N}) \delta_{c_{ij},0}, \quad c = \mathcal{O}(N^0)$
- indep random bonds J_{ij}
- disorder: $\{c_{ij}, J_{ij}\}$

Replica theory order parameters

Dependence on connectivity c
(average number of bonds/spin)

<i>connectivity</i>	<i>variables</i>	<i>order param</i>	<i>RS ansatz</i>
$c = N$	discrete	$\{q_{\alpha\beta}\}$	numbers, e.g. q
$c = N$	continuous	$\{q_{\alpha\beta}\}$	numbers, e.g. q
$1 \ll c \ll N$	discrete	$\{q_{\alpha\beta}\}$	numbers, e.g. q
$1 \ll c \ll N$	continuous	$\{q_{\alpha\beta}\}$	numbers, e.g. q
$c = \mathcal{O}(1)$	discrete	$P(\sigma_1, \dots, \sigma_n)$	functions, $P(h)$
$c = \mathcal{O}(1)$	continuous	$P(\sigma_1, \dots, \sigma_n)$	functionals, $W[\{P\}]$

Finite Connectivity Attractor Neural Networks

$$H = - \sum_{i < j}^N J_{ij} \sigma_i \sigma_j \quad J_{ij} = \frac{c_{ij}}{c} \phi \left(\sum_{\mu=1}^p \xi_i^\mu \xi_j^\mu \right)$$

$$P(c_{ij}) = \frac{c}{N} \delta_{c_{ij},1} + \left[1 - \frac{c}{N} \right] \delta_{c_{ij},0}$$

$$\sigma_i, \xi_i^\mu \in \{-1, 1\}$$

replica solution:

$$\begin{aligned} \bar{f} &= - \lim_{N \rightarrow \infty} (\beta N)^{-1} \overline{\log Z} \\ &= - \lim_{N \rightarrow \infty} \lim_{n \rightarrow 0} \frac{1}{\beta N n} \log \sum_{\boldsymbol{\sigma}^1 \dots \boldsymbol{\sigma}^n} e^{\frac{c}{2N} \sum_{ij} \left[e^{\frac{\beta}{c} \phi(\boldsymbol{\xi}_i \cdot \boldsymbol{\xi}_j) \sum_{\alpha} \sigma_i^\alpha \sigma_j^\alpha} - 1 \right]} \end{aligned}$$

analysis via sub-lattices:

$$I_{\boldsymbol{\xi}} = \{i \mid \xi_i = \boldsymbol{\xi}\}$$

$$\bar{f} = \lim_{n \rightarrow 0} \text{extr}_{\{P_{\boldsymbol{\xi}}(\boldsymbol{\sigma})\}} f[\{P_{\boldsymbol{\xi}}(\boldsymbol{\sigma})\}]$$

$$P_{\boldsymbol{\xi}}(\sigma_1, \dots, \sigma_n) = \lim_{N \rightarrow \infty} \frac{1}{|I_{\boldsymbol{\xi}}|} \sum_{i \in I_{\boldsymbol{\xi}}} \overline{\langle \prod_{\alpha} \delta_{\sigma_{\alpha}, \sigma_i^{\alpha}} \rangle}$$

Replica Symmetric ansatz ?

disorder: $D = \{c_{ij}, J_{ij}\}$

probability: $\mathcal{P}(D)$

$$\begin{aligned} P_{\xi}(\sigma_1, \dots, \sigma_n) &= \lim_{N \rightarrow \infty} \frac{1}{|I_{\xi}|} \sum_{i \in I_{\xi}} \overline{\langle \prod_{\alpha} \delta_{\sigma_{\alpha}, \sigma_i^{\alpha}} \rangle} \\ &= \lim_{N \rightarrow \infty} \frac{1}{|I_{\xi}|} \sum_{i \in I_{\xi}} \int dD \mathcal{P}(D) \langle \prod_{\alpha} \delta_{\sigma_{\alpha}, \sigma_i^{\alpha}} \rangle \Big|_D \end{aligned}$$

Ergodicity:

$$\langle \prod_{\alpha} \delta_{\sigma_{\alpha}, \sigma_i^{\alpha}} \rangle \Big|_D = \prod_{\alpha} p(\sigma_{\alpha} | D, i) = \prod_{\alpha} \left\{ \frac{e^{\beta h(D, i) \sigma_{\alpha}}}{2 \cosh[\beta h(D, i)]} \right\}$$

so

$$P_{\xi}(\sigma_1, \dots, \sigma_n) = \int dh W_{\xi}(h) \frac{e^{\beta h \sum_{\alpha} \sigma_{\alpha}}}{[2 \cosh(\beta h)]^n}$$

$$W_{\xi}(h) = \lim_{N \rightarrow \infty} \frac{1}{|I_{\xi}|} \sum_{i \in I_{\xi}} \int dD \mathcal{P}(D) \delta[h - h(D, i)]$$

RS order parameters:

2^p effective field distributions $W_{\xi}(h)$

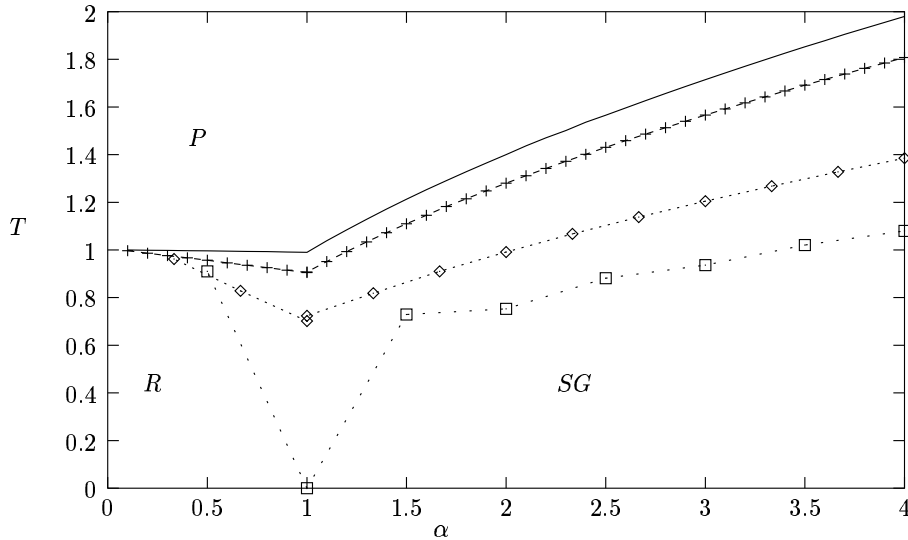
RS theory

$$W_{\xi}(h) = \int \frac{dm}{2\pi} e^{-imh} \times \exp \left\{ c \left\langle \int dh' W_{\xi'}(h') \left[e^{\frac{im}{\beta} \tanh^{-1}[\tanh(\beta h') \tanh(\frac{\beta}{c} \phi(\xi \cdot \xi'))]} - 1 \right] \right\rangle_{\xi'} \right\}$$

paramagnetic soln: $W_{\xi}(h) = \delta(h) \quad \forall \xi$

$$P \rightarrow R : \quad \frac{c}{p} 2^{-p} \sum_{n=0}^p \binom{p}{n} (p-2n) \tanh\left[\frac{\beta}{c} \phi(p-2n)\right] = 1$$

$$P \rightarrow SG : \quad c 2^{-p} \sum_{n=0}^p \binom{p}{n} \tanh^2\left[\frac{\beta}{c} \phi(p-2n)\right] = 1$$

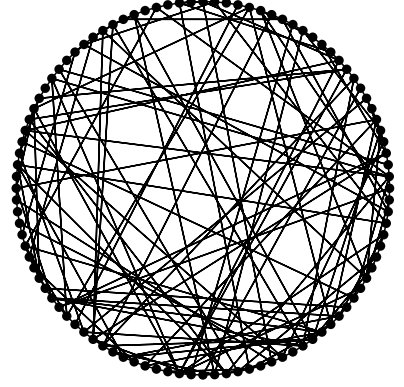


Fin conn Hopfield model: $\phi(x) = x$, $\alpha = p/c$
 $c = 2$ (\square), $c = 3$ (\diamond), $c = 10$ ($+$), $c = 100$ (line)

'Small World' Spin Models

N Ising spins on ring + random graph

$$H = - \sum_i J_i^s \sigma_i \sigma_{i+1} - \frac{1}{c} \sum_{i < j} J_{ij}^\ell c_{ij} \sigma_i \sigma_j$$



$N = 100, c = 2$

- $P(c_{ij}) = \frac{c}{N} \delta_{c_{ij},1} + (1 - \frac{c}{N}) \delta_{c_{ij},0}$
- indep random interaction energies $\{J_i^s, J_{ij}^\ell\}$
- disorder: $\{c_{ij}, J_{ij}\}$

Replica analysis

$$\bar{f} = \lim_{n \rightarrow 0} \frac{1}{n\beta} \text{extr}_{\{P\}} \left\{ \frac{c}{2} \sum_{\boldsymbol{\sigma} \boldsymbol{\sigma}'} P(\boldsymbol{\sigma}) P(\boldsymbol{\sigma}') \langle e^{\frac{\beta J_\ell}{c} \sum_\alpha \sigma_\alpha \sigma'_\alpha} - 1 \rangle_{J_\ell} - \log \lambda_0(n; P) \right\}$$

$$P(\boldsymbol{\sigma}) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_i \overline{\langle \prod_\alpha \delta_{\sigma_\alpha, \sigma_i^\alpha} \rangle}, \quad \boldsymbol{\sigma} = (\sigma_1, \dots, \sigma_n)$$

$\lambda_0(n; P)$: largest eigenvalue of $2^n \times 2^n$
replicated transfer matrix

$$\mathbf{T}(\boldsymbol{\sigma}, \boldsymbol{\sigma}'; P) = \langle e^{\beta J_s \sum_\alpha \sigma_\alpha \sigma'_\alpha + c \sum \mathbf{s} P(\mathbf{s}) \langle e^{\frac{\beta J_\ell}{c} \sum_\alpha \sigma_\alpha \sigma'_\alpha} - 1 \rangle_{J_\ell}} \rangle_{J_s}$$

RS theory

$$\mathbf{T}(\boldsymbol{\sigma}, \boldsymbol{\sigma}'; P) = \langle e^{\beta J_s \sum_{\alpha} \sigma_{\alpha} \sigma'_{\alpha} + c \sum_{\mathbf{s}} P(\mathbf{s}) \langle e^{\frac{\beta J_{\ell}}{c} \sum_{\alpha} \sigma_{\alpha} s_{\alpha} - 1 \rangle_{J_{\ell}}} \rangle_{J_s}$$

$$P(\boldsymbol{\sigma}) = \int dh W(h) \frac{e^{\beta h \sum_{\alpha=1}^n \sigma_{\alpha}}}{[2 \cosh(\beta h)]^n}$$

left/right

eigenvectors of \mathbf{T} :

$$\begin{aligned} v_L(\boldsymbol{\sigma}) &= \int dy \Psi(y) e^{\beta y \sum_{\alpha} \sigma_{\alpha}} \\ v_R(\boldsymbol{\sigma}) &= \int dx \Phi(x) e^{\beta x \sum_{\alpha} \sigma_{\alpha}} \end{aligned}$$

Final result, $n \rightarrow 0$:

$$\begin{aligned} \Phi(x) &= \int dx' \Phi(x') \sum_{k \geq 0} \frac{e^{-c} c^k}{k!} \int \prod_{r \leq k} dJ_{lr} dh_r P(J_{lr}) W(h_r) \\ &\quad \langle \delta[x - \sum_r A(J_{lr}/c, h_r) - A(J_s, x')] \rangle_{J_s} \end{aligned}$$

$$\begin{aligned} \Psi(y) &= \int dy' \Psi(y') \sum_{k \geq 0} \frac{e^{-c} c^k}{k!} \int \prod_{r \leq k} dJ_{lr} dh_r P(J_{lr}) W(h_r) \\ &\quad \langle \delta[y - A(J_s, y' + \sum_r A(J_{lr}/c, h_r))] \rangle_{J_s} \end{aligned}$$

$$W(h) = \int dx dy \Phi(x) \Psi(y) \delta[h - x - y]$$

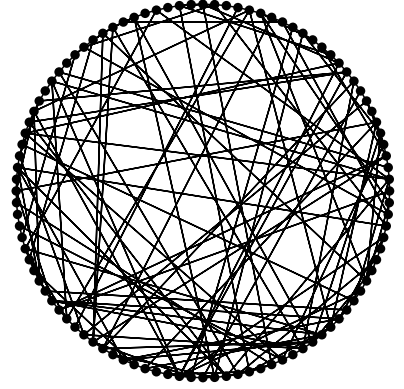
with

$$\begin{aligned} A(u, z) &= \frac{1}{2\beta} \log \left[\frac{\cosh(\beta(u+z))}{\cosh(\beta(u-z))} \right] \\ B(u, z) &= \frac{1}{2\beta} \log [4 \cosh(\beta(u+z)) \cosh(\beta(u-z))] \end{aligned}$$

Example: ‘Small World’ Magnets

$$H = -J_0 \sum_i \sigma_i \sigma_{i+1} - \frac{J}{c} \sum_{i < j} c_{ij} \sigma_i \sigma_j$$

(connectivity disorder only)



$$\begin{aligned} \Phi(x) = \int dx' \Phi(x') \sum_{k \geq 0} \frac{e^{-c} c^k}{k!} \int \prod_{r \leq k} [dh_r W(h_r)] \\ \times \delta[x - \sum_r A(\frac{J}{c}, h_r) - A(J_0, x')] \end{aligned}$$

$$\begin{aligned} \Psi(y) = \int dy' \Psi(y') \sum_{k \geq 0} \frac{e^{-c} c^k}{k!} \int \prod_{r \leq k} [dh_r W(h_r)] \\ \times \delta[y - A(J_0, y' + \sum_r A(\frac{J}{c}, h_r))] \end{aligned}$$

$$W(h) = \int dx dy \Phi(x) \Psi(y) \delta[h - x - y]$$

paramagnetic soln:

$$W(h) = \Phi(h) = \Psi(h) = \delta(h)$$

2nd order transitions away from P:

expand in moments of $W(h)$ and identify instabilities

$$\begin{aligned} \text{P} \rightarrow \text{F} : \quad & 1 = c \tanh\left(\frac{\beta J}{c}\right) \left[\frac{1 + \tanh(\beta J_0)}{1 - \tanh(\beta J_0)} \right] \\ \text{P} \rightarrow \text{SG} : \quad & \text{does not happen} \end{aligned}$$

$$P \rightarrow F : \quad \beta J = \frac{1}{2}c \log \left[\frac{c + e^{-2\beta J_0}}{c - e^{-2\beta J_0}} \right]$$

limits:

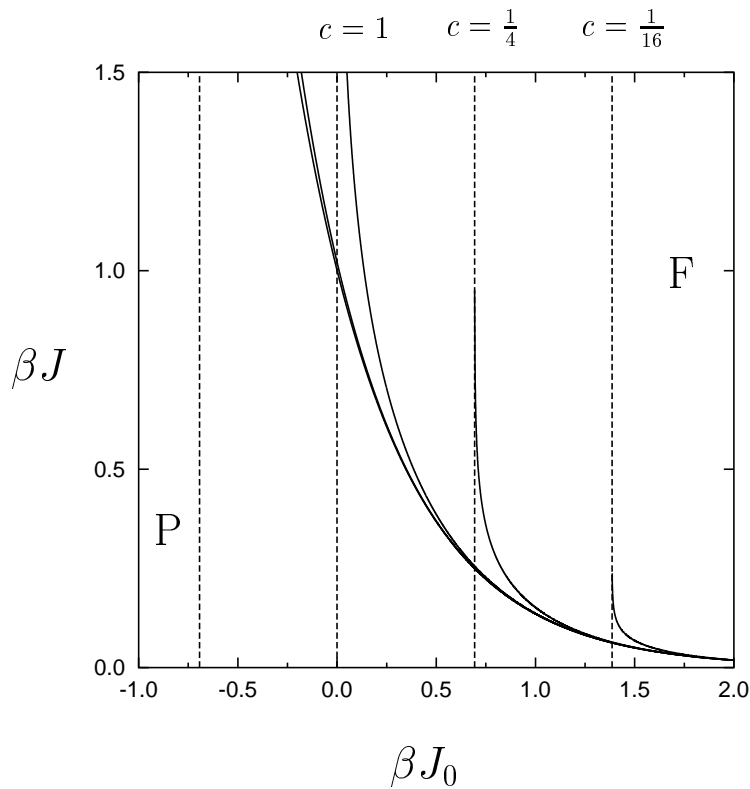
$$c \rightarrow \infty : \quad \beta J = e^{-2\beta J_0}$$

$$J_0 = 0 : \quad \text{order possible for } c > 1$$

$$J = 0 \text{ or } c = 0 : \quad \text{never order}$$

combination:

$$c > 0, J_0 > 0 : \quad \text{order for } T < T_c \text{ with } T_c(J_0, J, c) > 0 \\ \text{even for } c \ll 1 !!$$



solid:

P→F transition

for $c = \frac{1}{16}, \frac{1}{4}, 1, 4, 16$

dashed:

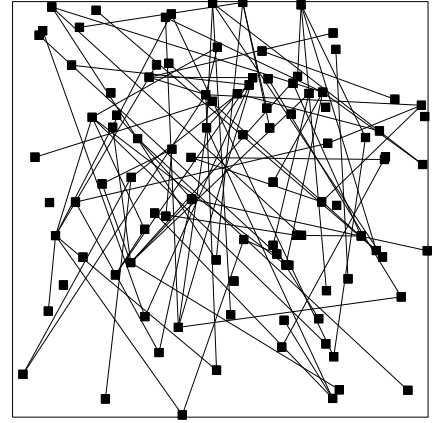
$\beta J_0 > \log(1/\sqrt{c})$

for $c = \frac{1}{16}, \frac{1}{4}, 1, 4, 16$

Soft Spin Models

N soft spins $\sigma_i \in \mathbb{R}$

$$H = - \sum_{i < j} c_{ij} J_{ij} \sigma_i \sigma_j + \sum_i V(\sigma_i)$$



- $P(c_{ij}) = \frac{c}{N} \delta_{c_{ij},1} + (1 - \frac{c}{N}) \delta_{c_{ij},0}$, $c = \mathcal{O}(N^0)$
- indep random bonds J_{ij}

Replica theory

order parameter:

$$P(\sigma_1, \dots, \sigma_n) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_i \overline{\langle \prod_{\alpha} \delta[\sigma_{\alpha} - \sigma_i^{\alpha}] \rangle}$$

saddle-point:

$$P(\boldsymbol{\sigma}) = \frac{e^{c \int d\boldsymbol{\sigma}' P(\boldsymbol{\sigma}') [\int dJ P(J) e^{\beta J \boldsymbol{\sigma} \cdot \boldsymbol{\sigma}' - 1] - \beta \sum_{\alpha} V(\sigma_{\alpha})}}{\int d\boldsymbol{\sigma}' e^{c \int d\boldsymbol{\sigma}'' P(\boldsymbol{\sigma}'') [\int dJ P(J) e^{\beta J \boldsymbol{\sigma}' \cdot \boldsymbol{\sigma}'' - 1] - \beta \sum_{\alpha} V(\sigma'_{\alpha})}}$$

Replica Symmetric ansatz ?

disorder: $D = \{c_{ij}, J_{ij}\}$
probability: $\mathcal{P}(D)$

$$\begin{aligned} P(\sigma_1, \dots, \sigma_n) &= \lim_{N \rightarrow \infty} \frac{1}{N} \sum_i \overline{\prod_{\alpha} \delta[\sigma_{\alpha} - \sigma_i^{\alpha}]} \\ &= \lim_{N \rightarrow \infty} \frac{1}{N} \sum_i \int dD \mathcal{P}(D) \left\langle \prod_{\alpha} \delta[\sigma_{\alpha} - \sigma_i^{\alpha}] \right\rangle \Big|_D \end{aligned}$$

Ergodicity:

$$\left\langle \prod_{\alpha} \delta[\sigma_{\alpha} - \sigma_i^{\alpha}] \right\rangle \Big|_D = \prod_{\alpha} P(\sigma_{\alpha} | D, i)$$

so

$$\begin{aligned} P_{\text{RS}}(\sigma_1, \dots, \sigma_n) &= \int \{dP\} W[\{P\}] \prod_{\alpha} P(\sigma_{\alpha}) \\ W[\{P\}] &= \lim_{N \rightarrow \infty} \frac{1}{N} \sum_i \int dD \mathcal{P}(D) \prod_{\sigma} \delta \left[P(\sigma) - \prod_{\alpha} P(\sigma | D, i) \right] \end{aligned}$$

RS order parameter for soft spins:
functional $W[\{P\}]$ on space of single-spin distributions $P(\sigma)$

saddle-point eqn, $n \rightarrow 0$:

$$\begin{aligned} W[\{P\}] &= \sum_{\ell \geq 0} \frac{c^{\ell}}{\ell!} e^{-c} \int \prod_{k \leq \ell} [\{dP_k\} W[\{P_k\}] dJ_k P(J_k)] \\ &\times \prod_{\sigma} \delta \left[P(\sigma) - \frac{e^{-\beta V(\sigma)} \prod_{k=1}^{\ell} \int d\sigma' P_k(\sigma') e^{\beta J_k \sigma \sigma'}}{\int d\sigma'' e^{-\beta V(\sigma'')} \prod_{k=1}^{\ell} \int d\sigma' P_k(\sigma') e^{\beta J_k \sigma'' \sigma'}} \right] \end{aligned}$$

Application: Chiral Soft Spin Models

N interacting vectorial soft spins

$\boldsymbol{\sigma}_i \in S_d$ (unit sphere in \mathbb{R}^d)

$$H = -J \sum_{i < j} c_{ij} \boldsymbol{\sigma}_i \cdot \mathbf{U}_{ij} \boldsymbol{\sigma}_j + \sum_i V(\boldsymbol{\sigma}_i)$$

- $P(c_{ij}) = \frac{c}{N} \delta_{c_{ij},1} + (1 - \frac{c}{N}) \delta_{c_{ij},0}$
- indep random unitary matrices \mathbf{U}_{ij} ,
 $P(\mathbf{U}) = P(\mathbf{U}^\dagger)$

$$W[\{P\}] = \sum_{\ell \geq 0} \frac{c^\ell}{\ell!} e^{-c} \int \prod_{k \leq \ell} [dP_k] W[\{P_k\}] d\mathbf{U}_k P(\mathbf{U}_k)]$$

$$\times \prod_{\boldsymbol{\sigma} \in S_d} \delta \left[P(\boldsymbol{\sigma}) - \frac{e^{-\beta V(\boldsymbol{\sigma})} \prod_{k=1}^{\ell} \int d\boldsymbol{\sigma}' P_k(\boldsymbol{\sigma}') e^{\beta J \boldsymbol{\sigma} \cdot \mathbf{U}_k \boldsymbol{\sigma}'}}{\int d\boldsymbol{\sigma}'' e^{-\beta V(\boldsymbol{\sigma}'')} \prod_{k=1}^{\ell} \int d\boldsymbol{\sigma}' P_k(\boldsymbol{\sigma}') e^{\beta J \boldsymbol{\sigma}'' \cdot \mathbf{U}_k \boldsymbol{\sigma}'}} \right]$$

d=2: Finitely Connected XY Spins with Random Chiral Interactions

$$\begin{aligned} \boldsymbol{\sigma}_i &\rightarrow (\cos \phi_i, \sin \phi_i), \\ P(\mathbf{U}) &\rightarrow P(\omega): \end{aligned} \quad H = -J \sum_{i < j} c_{ij} \cos(\phi_i - \phi_j - \omega_{ij})$$

$$\begin{aligned} W[\{P\}] &= \sum_{\ell \geq 0} \frac{c^\ell}{\ell!} e^{-c} \int \prod_{k \leq \ell} [\{dP_k\} W[\{P_k\}] d\omega_k P(\omega_k)] \\ &\times \prod_{\phi \in [0, 2\pi]} \delta \left[P(\phi) - \frac{\prod_{k=1}^{\ell} \int d\phi' P_k(\phi') e^{\beta J \cos(\phi - \phi' - \omega_k)}}{\int d\phi'' \prod_{k=1}^{\ell} \int d\phi' P_k(\phi') e^{\beta J \cos(\phi'' - \phi' - \omega_k)}} \right] \end{aligned}$$

paramagnetic soln:

$$W[\{P\}] = \prod_{\phi \in [0, 2\pi]} \delta \left[P(\phi) - \frac{1}{2\pi} \right]$$

2nd order transitions away from P:

- transform $P(\phi) \rightarrow \frac{1}{2\pi} + \Delta(\phi)$, with $W[\{P\}] \rightarrow \tilde{W}[\{\Delta\}]$
- expand in functional moments $\int \{d\Delta\} \tilde{W}[\{\Delta\}] \Delta(\phi_1) \dots \Delta(\phi_r)$
- identify instabilities

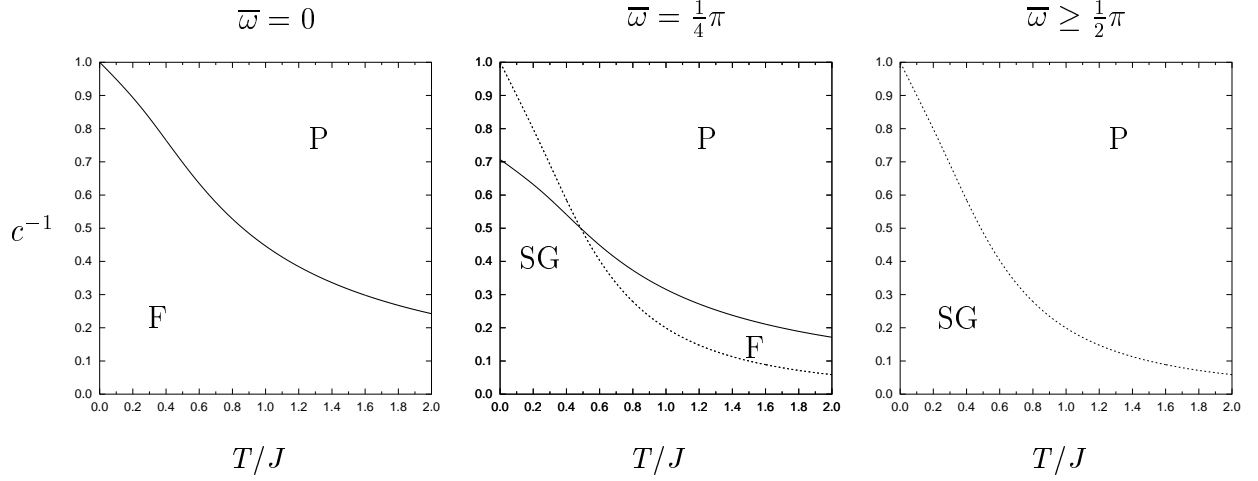
$$P \rightarrow F : \quad c^{-1} = \frac{I_1(\beta J)}{I_0(\beta J)} \int_{-\pi}^{\pi} d\omega P(\omega) \cos(\omega)$$

$$P \rightarrow SG : \quad c^{-1} = I_1^2(\beta J) / I_0^2(\beta J)$$

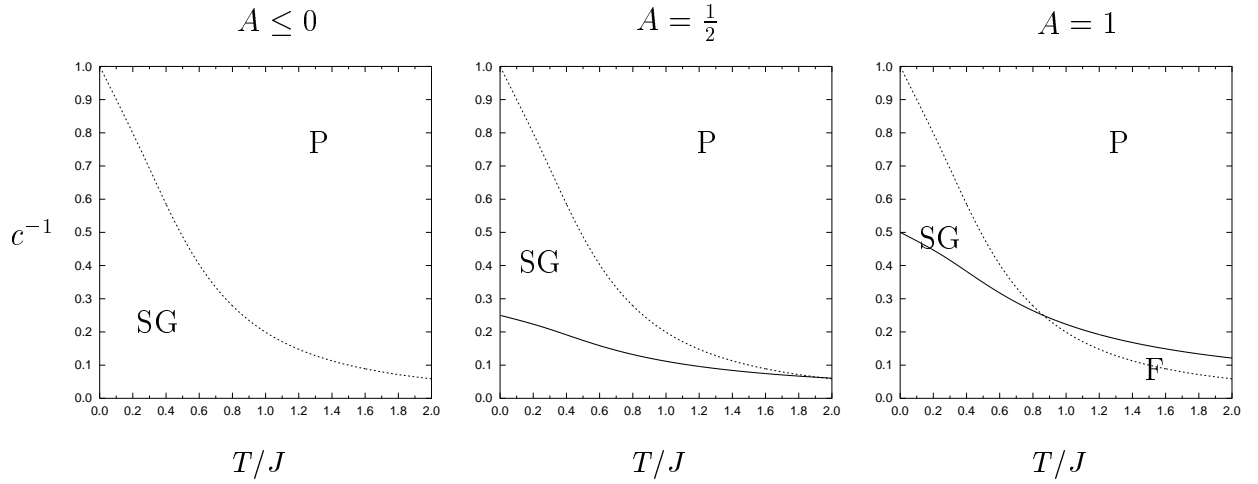
$$\text{modified Bessel func} \quad I_n(z) = \int_{-\pi}^{\pi} \frac{d\phi}{2\pi} \cos(n\phi) e^{z \cos(\phi)}$$

solid lines: P→F
dotted lines: P→SG

$$H = -J \sum_{i<j} c_{ij} \cos(\phi_i - \phi_j - \omega_{ij})$$



$$P(\omega) = \frac{1}{2}\delta(\omega - \bar{\omega}) + \frac{1}{2}\delta(\omega + \bar{\omega})$$



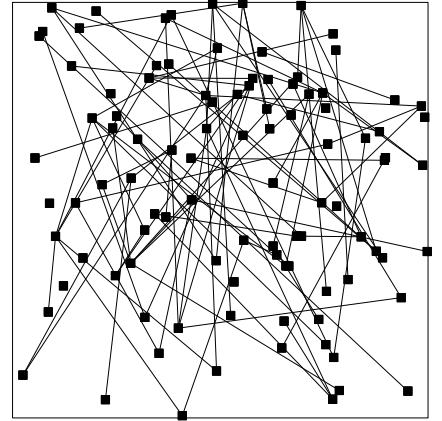
$$P(\omega) = \frac{1}{2\pi}[1 + A \cos(\ell\phi)]$$

Generating Functional Analysis For Finite Connectivity Models

finitely connected Ising spin model with
parallel stochastic dynamics:

$$p_{t+1}(\boldsymbol{\sigma}) = \sum_{\boldsymbol{\sigma}'} W_t[\boldsymbol{\sigma}; \boldsymbol{\sigma}'] p_t(\boldsymbol{\sigma}')$$

$$W_t[\boldsymbol{\sigma}; \boldsymbol{\sigma}'] = \prod_i \frac{e^{\beta \sigma_i h_i(\boldsymbol{\sigma}'; t)}}{2 \cosh[\beta h_i(\boldsymbol{\sigma}'; t)]}$$



local fields:

$$h_i(\boldsymbol{\sigma}; t) = \sum_{j \neq i} c_{ij} J_{ij} \sigma_j + \theta(t)$$

controlled symmetry:

$$i < j : \quad \text{Prob}(c_{ij}) = W(c_{ij})$$

$$i > j : \quad \text{Prob}(c_{ij}) = \epsilon_1 \delta_{c_{ij}, c_{ji}} + (1 - \epsilon_1) W(c_{ij})$$

$$i < j : \quad \text{Prob}(J_{ij}) = P(J_{ij})$$

$$i > j : \quad \text{Prob}(J_{ij}) = \epsilon_2 \delta[J_{ij} - J_{ji}] + (1 - \epsilon_2) P(J_{ij})$$

$$W(x) = \frac{c}{N} \delta_{x,1} + \left(1 - \frac{c}{N}\right) \delta_{x,0} \quad c = \mathcal{O}(N^0)$$

detailed balance &
equil stat mech:

$$\epsilon_1 = \epsilon_2 = 1$$

generating functional

$$\begin{aligned}
Z[\boldsymbol{\psi}] &= \overline{\langle e^{-i \sum_{i \leq N} \sum_{t < t_m} \psi_i(t) \sigma_i(t)} \rangle} \\
&= \sum_{\boldsymbol{\sigma}(0), \dots, \boldsymbol{\sigma}(t_m)} e^{-i \sum_i \sum_t \psi_i(t) \sigma_i(t)} p_0(\boldsymbol{\sigma}(0)) \overline{\prod_{t < t_m} W_t[\boldsymbol{\sigma}(t+1); \boldsymbol{\sigma}(t)]}
\end{aligned}$$

insert

$$1 = \prod_{it} \int dh_i(t) \delta[h_i(t) - \sum_{j \neq i} c_{ij} J_{ij} \sigma_j(t) - \theta(t)]$$

disorder average,

standard manipulations:

effective single spin problem

$P(\boldsymbol{\sigma}|\boldsymbol{\theta})$: fraction of sites i which exhibit
a single spin path $\boldsymbol{\sigma} = (\sigma(0), \sigma(1), \sigma(2), \dots)$
given a field path $\boldsymbol{\theta} = (\theta(0), \theta(1), \theta(2), \dots)$

$$P(\boldsymbol{\sigma}|\boldsymbol{\theta}) = p_0(\sigma(0)) e^{-c} \times$$

$$\left\{ \prod_t \left[\frac{e^{\beta \sigma(t+1) \theta(t)}}{2 \cosh[\beta \theta(t)]} \right] + \sum_{k > 0} \frac{c^k}{k!} \int dJ_1 P(J_1) \dots dJ_k P(J_k) \sum_{\boldsymbol{\sigma}'_1 \dots \boldsymbol{\sigma}'_k} \right. \\
\times \prod_{\ell=1}^k \left[(1 - \epsilon_1) P(\boldsymbol{\sigma}'_\ell | \mathbf{0}) + \epsilon_1 [\epsilon_2 P(\boldsymbol{\sigma}'_\ell | J_\ell \boldsymbol{\sigma}) + (1 - \epsilon_2) \langle P(\boldsymbol{\sigma}'_\ell | J' \boldsymbol{\sigma}) \rangle_{J'}] \right] \\
\left. \times \prod_t \frac{e^{\beta \sigma(t+1) [\theta(t) + \sum_{\ell \leq k} J_\ell \sigma'_\ell(t)]}}{2 \cosh[\beta [\theta(t) + \sum_{\ell \leq k} J_\ell \sigma'_\ell(t)]]} \right\}$$

$\epsilon_1 \in [0, 1]$: connect symm

$\epsilon_2 \in [0, 1]$: bond value symm

Simplest case: uniform bonds

$$P(J') = \delta(J - J')$$

$$P(\boldsymbol{\sigma}|\boldsymbol{\theta}) = p_0(\sigma(0)) \sum_{k \geq 0} \frac{e^{-c} c^k}{k!} \sum_{\boldsymbol{\sigma}'_1, \dots, \boldsymbol{\sigma}'_k} \prod_{\ell \leq k} [(1 - \epsilon_1) P(\boldsymbol{\sigma}'_\ell | \mathbf{0}) + \epsilon_1 P(\boldsymbol{\sigma}'_\ell | J\boldsymbol{\sigma})]$$

$$\times \prod_{t < t_m} \frac{e^{\beta \sigma(t+1)[\theta(t) + J \sum_{\ell \leq k} \sigma'_\ell(t)]}}{2 \cosh[\beta[\theta(t) + J \sum_{\ell \leq k} \sigma'_\ell(t)]]}$$

- asymmetric connectivity, $\epsilon_1 = 0$:

$$\text{let } W(\boldsymbol{\sigma}) \equiv P(\boldsymbol{\sigma} | \mathbf{0})$$

$$W(\boldsymbol{\sigma}) = p_0(\sigma(0)) \sum_{k \geq 0} \frac{e^{-c} c^k}{k!} \sum_{\boldsymbol{\sigma}'_1, \dots, \boldsymbol{\sigma}'_k} \prod_{\ell \leq k} W(\boldsymbol{\sigma}'_\ell)$$

$$\times \prod_{t < t_m} \frac{e^{\beta J \sigma(t+1) \sum_{\ell \leq k} \sigma'_\ell(t)}}{2 \cosh[\beta J \sum_{\ell \leq k} \sigma'_\ell(t)]}$$

(\rightarrow Derrida, Gardner & Zippelius)

- symmetric connectivity, $\epsilon_1 = 1$:

$$\text{let } W(\boldsymbol{\sigma}|\boldsymbol{\tau}) \equiv P(\boldsymbol{\sigma} | J\boldsymbol{\tau})$$

$$W(\boldsymbol{\sigma}|\boldsymbol{\tau}) = p_0(\sigma(0)) \sum_{k \geq 0} \frac{e^{-c} c^k}{k!} \sum_{\boldsymbol{\sigma}'_1, \dots, \boldsymbol{\sigma}'_k} \prod_{\ell \leq k} W(\boldsymbol{\sigma}'_\ell | \boldsymbol{\sigma})$$

$$\times \prod_{t < t_m} \frac{e^{\beta J \sigma(t+1)[\tau(t) + \sum_{\ell \leq k} \sigma'_\ell(t)]}}{2 \cosh[\beta J [\tau(t) + \sum_{\ell \leq k} \sigma'_\ell(t)]]}$$

OUTLOOK

statics

- low temperature solutions and F→SG transitions, beyond variational ansätze ?
- replica symmetry breaking, generalization of Parisi scheme, or something else ?
- ‘small world’ spin glasses, with RKKY-type bonds
nearest neighbour: $J_0 > 0$
random finite conn: J_{ij} zero average Gaussian

dynamics

- solutions of GFA order parameter equations
- exact dynamical replica theory ?
(single time observables rather than path statistics !)

other

- finite connectivity Gardner calculations,
(storage capacity of attractor networks)
- finite connectivity quantum systems,
e.g. $H = -J \sum_{i<j} c_{ij} \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$ (with Pauli matrices)